Conclusion

Study of the effect of parameters on the decay rate of a fourth order problem Annual Meeting of the COMPLEX Doctoral School 2023

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Monday 13 November 2023

We study the evolution equation

 $\partial_{tt}u(x,t) + a\partial_{xxxx}u(x,t) + b\partial_tu(x,t) + \alpha\partial_tu(\xi,t)\delta_{\xi} = 0,$ where $(x,t) \in (0,1) \times (0,+\infty).$

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- *u*: transverse displacement of the bridge deck (identified with [0, 1]);
- δ_{ξ} : presence of a shape memory alloy cable at $x = \xi$.

Conclusion

Boundary and initial conditions

We couple the equations with boundary conditions

$$u(0,t) = u(1,t) = \partial_x^2 u(0,t) = \partial_x^2 u(1,t) = 0, \qquad t \in (0,+\infty),$$

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 $t \in (0,+\infty),$

and initial conditions

 $u(x,0) = u_0(x), \quad \partial_t u(x,0) = u_1(x), \qquad x \in (0,1).$

We define the energy of a solution u(x, t) by

$$E(t) := \frac{1}{2} \int_0^1 \left(|\partial_t u|^2 + \mathbf{a} |\partial_{xx} u|^2 \right) \mathrm{d}x.$$

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Differentiating (formally) and integrating by parts, we obtain

$$\partial_t E(t) = -b\left(\int_0^1 |\partial_t u|^2 \,\mathrm{d}x\right) - \alpha |\partial_t u(\xi, t)|^2$$

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Hence, the system is dissipative, in the sense that the energy decreases.

The damping rate

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Theorem (Régnier (2022))

The optimal energy decay rate of the equation, i.e. the smallest $\omega(\alpha) < 0$ such that, for all initial conditions U_0 , there exists $C(U_0) > 0$ with

$$E(t) \leq C(U_0)e^{2\omega(\alpha)t}$$

for all $t \ge 0$ is given by

$$\omega(\alpha) := \sup \Big\{ \Re(\mu) \mid \mu \text{ is an eigenvalue of the problem} \Big\}.$$

Conclusion

The role of the parameter α

A natural question

How does the decay rate $\omega(\alpha)$ depend on α ?

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An important assumption

To avoid any "resonance phenomena", we will assume that $\xi \notin \mathbb{Q}$.

Proposition

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 $\begin{aligned} &(\mu+b)\sinh(\lambda)\sin(\lambda)\\ &+\alpha\lambda\Big[\sin(\lambda)\sinh(\lambda\xi)\sinh(\lambda(1-\xi))-\sinh(\lambda)\sin(\lambda\xi)\sin(\lambda(1-\xi))\Big]=0,\end{aligned}$

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$$\lambda(\mu) := \sqrt[4]{-rac{b\mu+\mu^2}{a}}.$$

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$$\lambda(\mu) := \sqrt[4]{-rac{b\mu+\mu^2}{a}}.$$

Remark

Replacing λ by $i\lambda$, $-\lambda$ or $-i\lambda$ leads to an equivalent equation.

Finding the eigenvalues of the problem amounts to find roots of the function

$$F(\mu; \boldsymbol{\alpha}) := 2\mu(\mu + b)F_0(\lambda) + \boldsymbol{\alpha}\mu\lambda F_1(\lambda),$$

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- $\alpha = 0 \rightsquigarrow roots of F_0;$
- $\alpha \to +\infty \rightsquigarrow$ roots of F_1 .

Dependence of the roots on parameters A general fact from complex analysis

Theorem ("Holomorphic implicit function Theorem", very roughly stated)

Roots of holomorphic functions depend **continuously**, **including multiplicities**, on the parameters, and the branches of roots are holomorphic.

Dependence of the roots on parameters

A simple example: roots of $z \mapsto z^2 + c$

$$z\mapsto z^2-4$$



Blue: values. Red: multiplicities.

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Dependence of the roots on parameters

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Dependence of the roots on parameters A simple control of the roots of parameters

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The case $\alpha = 0$: roots of $\lambda \mapsto F_0(\lambda)$ A computation

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Therefore, the set of roots of F_0 is

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and all have multiplicity one, except zero which has multiplicity two.
The case $\alpha = 0$: roots of $\lambda \mapsto F_0(\lambda)$

Graphical representation in the λ plane



Conclusion

The case $\alpha = 0$: roots of $\mu \mapsto F_0(\lambda(\mu))$ Graphical representation in the μ plane (a = 0.05, b = 3)

$$\lambda(\mu) = \sqrt[4]{-\frac{b\mu + \mu^2}{a}}$$

i.e.
$$\mu^2 + b\mu + a\lambda^4 = 0$$

so that
$$\mu(\lambda) = \frac{-b \pm \sqrt{b^2 - 4a\lambda^4}}{2}$$

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Note: 0 is a root, but is *not* an eigenvalue!

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Roots of $\lambda \mapsto F_1(\lambda)$ The strategy: a continuation argument

We write

$$F_1(\lambda) = s(\lambda) - t(\lambda)$$

where

$$s(\lambda) := \sin(\lambda) \sinh(\lambda\xi) \sinh(\lambda(1-\xi))$$

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Strategy: study roots of

$$\widetilde{F}_{\gamma}(\lambda) := s(\lambda) - \gamma t(\lambda).$$

as γ varies from 0 to 1.

Roots of $\lambda \mapsto F_1(\lambda)$ Roots of *s* and *t*: a computation

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$$\begin{split} s(\lambda) &:= \sin(\lambda) \sinh(\lambda\xi) \sinh(\lambda(1-\xi)) \\ t(\lambda) &:= \sinh(\lambda) \sin(\lambda\xi) \sin(\lambda(1-\xi)), \end{split}$$

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so that

$$\left\{\lambda \in \mathbb{C} \mid s(\lambda) = 0\right\} = \left\{k\pi, i\frac{k\pi}{\xi}, i\frac{k\pi}{1-\xi} \mid k \in \mathbb{Z}\right\}$$

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All those roots have multiplicity one, except 0.

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Introd	uction

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Roots of $\lambda \mapsto F_1(\lambda)$



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Theorem (Rayleigh (1894) - Beatty (1927))

Let 0 < r < 1 be irrational. Define the sets

$$A := \left\{ \left\lfloor \frac{n}{r} \right\rfloor \mid n \in \mathbb{Z}^{>0} \right\}, \qquad B := \left\{ \left\lfloor \frac{n}{1-r} \right\rfloor \mid n \in \mathbb{Z}^{>0} \right\}.$$

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- S. Beatty "Problem 3173". American Mathematical Monthly. 33 (3):
 p. 159 (1926).

Roots of $\lambda \mapsto F_1(\lambda)$ Beatty's Theorem: a numerical example

Let us take $r = \sqrt{2} - 1$. Then (using a little script),

$$A = \left\{ \left\lfloor \frac{n}{r} \right\rfloor \mid n \in \mathbb{Z}^{>0} \right\}$$
$$= \left\{ 2, 4, 7, 9, 12, 14, 16, 19, \dots \right\}$$

and

$$B = \left\{ \left\lfloor \frac{n}{1-r} \right\rfloor \mid n \in \mathbb{Z}^{>0} \right\}$$
$$= \left\{ 1, 3, 5, 6, 8, 10, 11, 13, 15, 17, 18, 20, \dots \right\}.$$

Study of the decay rate

Conclusion

Roots of $\lambda \mapsto F_1(\lambda)$



Study of the decay rate

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Study of the decay rate

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Roots of $\mu \mapsto F_1(\lambda(\mu))$ Graphical representation in the μ plane



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Moving α from 0 to $+\infty$

Graphical representation in the μ plane: $\textit{a} = 0.05, \textit{b} = 3, \xi = \sqrt{2} - 1$



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Pink: roots of F_0 ($\alpha = 0$)

Violet: roots of F_1 ($\alpha \to +\infty$)

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Graphical representation in the μ plane: $a = 0.05, b = 3, \xi = \sqrt{2} - 1$



Pink: roots of F_0 ($\alpha = 0$)

Violet: roots of $F_1 (\alpha \to +\infty)$

Orange: eigenvalues ($\alpha = 2$)

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Moving α from 0 to $+\infty$



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Theorem (G., Régnier, Troestler (2023))

Recall that the optimal decay rate $\omega(\alpha)$ is given by

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- **1** ω is continuous in α .
- **2** ω is nondecreasing;

3 one has

$$\lim_{\alpha \to +\infty} \omega(\alpha) = 0. \quad (!)$$

A physical conclusion

The term $\alpha \partial_t u(\xi, t) \delta_{\xi}$ is **definitely not** a damping term in the problem.

Thanks for your attention!

Thanks!

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